# Flat Vector Bundles and Open Covers 

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To the memory of B. Yu. Sternin


#### Abstract

We establish a generic counting formula for the Euler number of a flat vector bundle of rank $2 n$ over a $2 n$-dimensional closed manifold in terms of vertices of transversal open covers of the underlying manifold. We use the Mathai-Quillen formalism to prove our result.


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## 1. Introduction

By the celebrated Gauss-Bonnet-Chern theorem [1] and its ramifications, if a real vector bundle $F$ of rank $2 n$ over a $2 n$-dimensional closed smooth manifold $M$ admits a flat connection preserving a certain Euclidean metric on $F$, then the Euler number of $F$ vanishes, $\langle e(F),[M]\rangle=0$, where $e(F)$ is the Euler class of $F$. Milnor [2] constructs examples of rank two flat vector bundles over a closed surface of genus $g \geq 2$ such that the corresponding Euler number is nonzero. A closely related question is the famous Chern conjecture on the vanishing of the Euler characteristic of a closed affine manifold (cf. [3]).

Our original motivation is to understand the Chern conjecture by using the Mathai-Quillen formalism on the geometric construction of the Thom class [4]. While our efforts [5] on proving the Chern conjecture in this way have not been successful, we find that our method leads us to a kind of mysterious counting formula for the Euler number of a flat vector bundle in terms of transversal open covers of the underlying manifold. The purpose of this paper is to present this mysterious formula.

This paper is organized as follows. In Sec. 2, we present an exterior algebra version of the Mathai-Quillen formalism, which allows us to avoid the usual difficulty in computing the Euler class that the connection in question need not
preserve any metric on a given vector bundle. In Sec. 3, we prove our main result, which is stated as Theorem 3.4 in Sec. 3.1.

## 2. Exterior Algebra Version of the Mathai-Quillen Formalism

In this section, we present an exterior algebra version of the Mathai-Quillen formalism [4]. That is, we replace the spinor bundle considered in [4] by the exterior algebra bundle. This gives a formula for the Euler class for an arbitrary connection on a Euclidean vector bundle.

Let $\pi: E \rightarrow M$ be a real oriented vector bundle of even rank over a closed oriented manifold $M$. Let $\nabla^{E}$ be a connection on $E$. Then it induces a connection $\nabla^{\Lambda^{*}\left(E^{*}\right)}$ on $\Lambda^{*}\left(E^{*}\right)$ preserving the $\mathbf{Z}_{2}$-splitting $\Lambda^{*}\left(E^{*}\right)=\Lambda^{\text {even }}\left(E^{*}\right) \oplus \Lambda^{\text {odd }}\left(E^{*}\right)$.

Let $g^{E}$ be a Euclidean metric on $E$. For any $Z \in E$, let $Z^{*} \in E^{*}$ be its metric dual. As usual (cf. [6, Sec. 4.3]), let $c(Z)$ be the Clifford action on $\Lambda^{*}\left(E^{*}\right)$ defined by

$$
\begin{equation*}
c(Z)=Z^{*} \wedge-i_{Z} \tag{1}
\end{equation*}
$$

where $Z^{*} \wedge$ (respectively, $i_{Z}$ ) is the exterior (respectively, interior) multiplication by $Z^{*}$ (respectively, $Z$ ). Then

$$
\begin{equation*}
c(Z)^{2}=-|Z|_{g E}^{2} \tag{2}
\end{equation*}
$$

Let $A$ be the superconnection [7] on $\pi^{*} \Lambda^{*}\left(E^{*}\right)$ defined by

$$
\begin{equation*}
A=\pi^{*} \nabla^{\Lambda^{*}\left(E^{*}\right)}+c(Z) \tag{3}
\end{equation*}
$$

where $c(Z)$ now acts on $\left.\left(\pi^{*} \Lambda^{*}\left(E^{*}\right)\right)\right|_{Z}$ by

$$
c(Z)\left(\left.\left(\pi^{*} \alpha\right)\right|_{Z}\right)=\left.\left(\pi^{*}(c(Z) \alpha)\right)\right|_{Z}
$$

for any $\alpha \in \Lambda^{*}\left(E^{*}\right)$. By (2) and (3), $\exp \left(A^{2}\right)$ is of exponential decay along vertical directions in $E$.
Theorem 2.1. The closed form $\left(\frac{1}{2 \pi}\right)^{\mathrm{rk}(E)}\left\{\int_{E / M} \operatorname{tr}_{s}\left[\exp \left(A^{2}\right)\right]\right\}^{(\operatorname{rk}(E))}$ on $M$ is a representative of $e(E)$, where $e(E)$ is the Euler class of $E$.
Proof. By the Chern-Weil theory for superconnections (cf. [8, Prop. 1.43]) and the above-mentioned exponential decay property of $\exp \left(A^{2}\right)$, we see that the cohomology class represented by $\int_{E / M} \operatorname{tr}_{s}\left[\exp \left(A^{2}\right)\right]$ is independent of the choice of $\nabla^{E}$ and $g^{E}$. Thus, we may well assume that $\nabla^{E}$ preserves $g^{E}$. Then one can follow the strategy in [4].

In fact, since the computation is local, one may well assume that $E$ is spin. Then one has the following decomposition (cf. [9]) in terms of the (Hermitian) spinor bundle $S(E)=S_{+}(E) \oplus S_{-}(E)$ associated with $\left(E, g^{E}\right)$ :

$$
\begin{equation*}
\Lambda^{*}\left(E^{*}\right)=\left(S_{+}(E) \oplus S_{-}(E)\right) \widehat{\otimes}\left(S_{+}^{*}(E) \oplus S_{-}^{*}(E)\right) \tag{4}
\end{equation*}
$$

and $c(Z)$ now acts on $\left.\left(\pi^{*} S(E)\right)\right|_{Z}$. Moreover, $\nabla^{\Lambda^{*}\left(E^{*}\right)}$ has the decomposition into $\nabla^{S(E)} \otimes \operatorname{Id}_{S^{*}(E)}+\operatorname{Id}_{S(E)} \otimes \nabla^{S^{*}(E)}$, where $\nabla^{S(E)}$ and $\nabla^{S^{*}(E)}$ are the induced Hermitian connections on $S(E)$ and $S^{*}(E)$, respectively.

From (3) and (4), one obtains

$$
\begin{equation*}
\operatorname{tr}_{s}\left[\exp \left(A^{2}\right)\right]=\operatorname{tr}_{s}\left[\exp \left(\left(\pi^{*} \nabla^{S(E)}+c(Z)\right)^{2}\right)\right] \cdot \operatorname{tr}_{s}\left[\exp \left(\left(\pi^{*} \nabla^{S^{*}(E)}\right)^{2}\right)\right] \tag{5}
\end{equation*}
$$

By [4, Thm. 4.5], one has

$$
\left(\frac{\sqrt{-1}}{2 \pi}\right)^{\frac{\mathrm{rk}(E)}{2}} \operatorname{tr}_{s}\left[\exp \left(\left(\pi^{*} \nabla^{S(E)}+c(Z)\right)^{2}\right)\right]=(-1)^{\frac{\mathrm{rk}(E)}{2}} \operatorname{det}\left(\frac{\sinh \left(\pi^{*} R^{E} / 2\right)}{\pi^{*} R^{E} / 2}\right)^{\frac{1}{2}} U
$$

where $R^{E}=\left(\nabla^{E}\right)^{2}$ and $U$ is the Thom form constructed in [4, (4.7)].
By [9, Prop. III.11.24], one has

$$
\begin{equation*}
\operatorname{ch}\left(S_{+}^{*}(E)-S_{-}^{*}(E)\right)=\frac{e(E)}{\widehat{A}(E)} . \tag{6}
\end{equation*}
$$

By (5)-(6) and [4, Thm. 4.10], which integrates $U$ along vertical fibers, one sees that $\left(\frac{1}{2 \pi}\right)^{\mathrm{rk}(E)}\left\{\int_{E / M} \operatorname{tr}_{s}\left[\exp \left(A^{2}\right)\right]\right\}^{(\mathrm{rk}(E))}$ is a representative of $e(E)$.

Corollary 2.2. If $\operatorname{rk}(E)=\operatorname{dim} M$, then

$$
\langle e(E),[M]\rangle=\left(\frac{1}{2 \pi}\right)^{\mathrm{rk}(E)} \int_{E} \operatorname{tr}_{s}\left[\exp \left(A^{2}\right)\right] .
$$

## 3. Counting Formula for the Euler Number of Flat Vector Bundles

In this section, we state and prove a generic counting formula for the Euler number of a flat vector bundle in terms of transversal open covers.

This section is organized as follows. In Sec. 3.1, we present the basic setting and state our main result as Theorem 3.4. In Sec. 3.2, we present an application of Corollary 2.2 to the special case of flat vector bundles. In Secs. 3.3-3.5, we prove Theorem 3.4.

### 3.1. Flat Vector Bundles and the Counting Formula

Let $\pi: F \rightarrow M$ be a real oriented flat vector bundle of rank $2 n$ over a $2 n$-dimensional closed oriented manifold $M$. Let $\nabla^{F}$ denote the underlying flat connection on $F$. Then there exists a finite collection of open coordinate charts $\left\{\left(U_{\alpha},\left(x_{\alpha}^{i}\right)\right)\right\}$, $\alpha=1, \cdots, N,{ }^{1}$ covering $M$ such that $\nabla^{F}$ induces a canonical identification

$$
\left.F\right|_{U_{\alpha}} \simeq U_{\alpha} \times \mathbf{R}_{\alpha}^{2 n}
$$

over each $U_{\alpha}$. For any $U_{\alpha}$, we fix a coordinate system $\left(y_{\alpha}^{i}\right)$ of $\mathbf{R}_{\alpha}^{2 n}$. Then $\left(y_{\alpha}^{i}\right)$ and $\left(y_{\beta}^{i}\right)$ are related to each other over $U_{\alpha} \cap U_{\beta}$ by a constant linear transformation. Moreover, the horizontal exterior differential and the vertical exterior differential

$$
d^{H}=\sum_{i} d x_{\alpha}^{i} \frac{\partial}{\partial x_{\alpha}^{i}}, \quad d^{V}=\sum_{i} d y_{\alpha}^{i} \frac{\partial}{\partial y_{\alpha}^{i}}
$$

[^0]on $\left.F\right|_{U_{\alpha}}$ are independent of $\alpha$, the exterior differential on the total space $\mathcal{F}$ of $F$ has the decomposition
\[

$$
\begin{equation*}
d=d^{H}+d^{V}, \tag{7}
\end{equation*}
$$

\]

and

$$
\left(d^{H}\right)^{2}=\left(d^{V}\right)^{2}=d^{H} d^{V}+d^{V} d^{H}=0 .
$$

Without loss of generality, we assume that each $U_{\alpha}$ has smooth boundary $\partial U_{\alpha}$ and that these boundaries intersect each other completely transversally. Such an open cover of $M$ will be called a transversal open cover. ${ }^{2}$ Then each point of $M$ lies on at most $2 n$ distinct boundaries. Moreover, the set

$$
\mathbf{B}=\{p \in M: p \text { lies on } 2 n \text { distinct boundaries }\}
$$

is finite.
Let $h:[0,1] \rightarrow[0,1]$ be the smooth function $h(t)=\exp \left(-1 / t^{2}\right)$.
Let $g^{T M}$ be any metric on $T M$. For each $U_{\alpha}$, let $r_{\alpha}$ be the normal geodesic coordinate near $\partial U_{\alpha}$. Let $\rho_{\alpha} \in C^{\infty}(M)$ be a function such that

$$
\begin{equation*}
\rho_{\alpha}=h\left(r_{\alpha}\right) \tag{8}
\end{equation*}
$$

in $U_{\alpha}$ near $\partial U_{\alpha}, \operatorname{Supp}\left(\rho_{\alpha}\right) \subseteq \bar{U}_{\alpha}$, and

$$
\begin{equation*}
\rho_{\alpha}>0 \quad \text { in } U_{\alpha} . \tag{9}
\end{equation*}
$$

The existence of $\rho_{\alpha}$ is clear.
For any function or a smooth form $\rho$ on $M$, we use the same notation $\rho$ to denote its lift $\pi^{*} \rho$.

For any $p \in \mathbf{B}$, set $U_{p}=\bigcap_{p \in U_{\alpha}} U_{\alpha}$. Then there exists a (sufficiently small) open neighborhood $W_{p}$ of $p \in M$ with $\bar{W}_{p} \subset U_{p}$ such that $\bar{W}_{p} \cap \bar{W}_{q}=\emptyset$ for any distinct $p, q \in \mathbf{B}$.

Remark 3.1. For any $p \in \mathbf{B}$, let $U_{\alpha_{1}}, \ldots, U_{\alpha_{N_{p}}}, \alpha_{1}>\cdots>\alpha_{N_{p}}$, be the open coordinate charts containing $p$. Without loss of generality, we assume that $\rho_{\alpha_{1}}=$ $\cdots=\rho_{\alpha_{N_{p}}}=1$ in $W_{p}$. Let $U_{\beta_{1}}, \ldots, U_{\beta_{2 n}}, \beta_{1}>\cdots>\beta_{2 n}$, be the open coordinate charts such that $\partial U_{\beta_{1}}, \ldots, \partial U_{\beta_{2 n}}$ meet at $p$. Moreover, by taking $W_{p}$ to be sufficiently small, we may assume that the intersection $\bar{W}_{p} \cap \bar{U}_{\alpha}$ is empty for every $\alpha \notin\left\{\alpha_{1}, \ldots, \alpha_{N_{p}}, \beta_{1}, \ldots, \beta_{2 n}\right\}$.

As a final notation, we set

$$
\mathbf{B}_{+}=\left\{p \in \mathbf{B}: \beta_{2 n}>\alpha_{1}\right\} .
$$

The following proposition will be proved in Sec. 3.4.

[^1]Proposition 3.2. For any $p \in \mathbf{B}_{+}$, there exists a sufficiently small open neighborhood $V_{p} \subset W_{p}$ of $p$ such that the limit

$$
\begin{align*}
\nu_{p}=\lim _{T \rightarrow+\infty} \int_{\mathcal{F}} \phi \prod_{i=1}^{2 n}\left(\frac{d \rho_{\beta_{i}} T^{\beta_{i}}}{4 \pi}\right. & \left.\sum_{j=1}^{2 n} d\left(y_{\beta_{i}}^{j}\right)^{2}\right) \\
& \cdot \exp \left(-\sum_{i, j=1}^{2 n} \rho_{\beta_{i}} T^{\beta_{i}}\left(y_{\beta_{i}}^{j}\right)^{2}-T^{\alpha_{1}} \sum_{j=1}^{2 n}\left(y_{\alpha_{1}}^{j}\right)^{2}\right) \tag{10}
\end{align*}
$$

exists for any $\phi \in C^{\infty}(M)$ supported in $V_{p}$ with $\phi=1$ near $p$.
Remark 3.3. The local index $\nu_{p}$ in (10) essentially depends only on the correlations between $\sum_{j=1}^{2 n}\left(y_{\alpha_{1}}^{j}\right)^{2}$ and $\sum_{j=1}^{2 n}\left(y_{\beta_{i}}^{j}\right)^{2}, i=1, \ldots, 2 n$. For example, if there exist constants $a_{i}, i=1, \ldots, 2 n$, such that

$$
\sum_{i, j=1}^{2 n} a_{i}\left(y_{\beta_{i}}^{j}\right)^{2}=\sum_{j=1}^{2 n}\left(y_{\alpha_{1}}^{j}\right)^{2}
$$

on $\pi^{-1}(p)$, then $\nu_{p}=0$. This fact might be of help when studying affine manifolds. Now we can state our main result as follows.

Theorem 3.4. The following identity holds:

$$
\begin{equation*}
\langle e(F),[M]\rangle=\sum_{p \in \mathbf{B}_{+}} \nu_{p} . \tag{11}
\end{equation*}
$$

Theorem 3.4 will be proved in Sec. 3.5.
Remark 3.5. Milnor's above-mentioned example shows that the index $\nu_{p}$ can be nonzero. On the other hand, the sum on the right-hand side in (11) looks mysterious. While it should be related to Čech cohomology (at least in the case of $F=T M$ ), it depends on the ordering of the coordinate charts $U_{\alpha}$. If one changes the ordering, then the set $\mathbf{B}_{+}$changes. This sounds interesting and deserves further study.

### 3.2. Superconnections and Flat Vector Bundles

In what follows, for clarity, we use the hat $\hat{\cdot}$ over symbols to indicate elements of $\pi^{*} \Lambda^{*}\left(F^{*}\right)$. Set

$$
\begin{equation*}
\widehat{Y}=\sum_{k=1}^{2 n} y_{\alpha}^{k} \frac{\widehat{\partial}}{\partial y_{\alpha}^{k}} \in \Gamma\left(\pi^{*} F\right) \tag{12}
\end{equation*}
$$

for each $\pi^{-1}\left(U_{\alpha}\right)$. Then $\widehat{Y}$ is a well-defined canonical section of $\pi^{*} F$ over the total space $\mathcal{F}$ of $F$.

For any $T>0$, let $\widehat{\eta}_{T} \in \Gamma\left(\pi^{*} F^{*}\right)$ be defined by

$$
\begin{equation*}
\widehat{\eta}_{T}=\sum_{\alpha=1}^{N} \rho_{\alpha} T^{\alpha} \sum_{k} y_{\alpha}^{k} \widehat{d y_{\alpha}^{k}} . \tag{13}
\end{equation*}
$$

Remark 3.6. For each $T>0$, if we equip $F$ with the Euclidean metric

$$
g_{T}^{F}=\sum_{\alpha=1}^{N} \rho_{\alpha} T^{\alpha} \sum_{k}\left(d y_{\alpha}^{k}\right)^{2}
$$

then $\widehat{\eta}_{T}$ is the metric dual of $\widehat{Y}$ (with respect to $\pi^{*} g_{T}^{F}$ ).
Set

$$
\begin{equation*}
c_{T}(\widehat{Y})=\widehat{\eta}_{T} \wedge-i_{\widehat{Y}}, \tag{14}
\end{equation*}
$$

as in (1). Then $c_{T}(\widehat{Y})$ acts on $\left.\left(\pi^{*} \Lambda^{*}\left(F^{*}\right)\right)\right|_{Y}$. Moreover,

$$
\begin{equation*}
|Y|_{g_{T}^{F}}^{2}=-c_{T}(\widehat{Y})^{2}=\sum_{\alpha, k} \rho_{\alpha} T^{\alpha}\left(y_{\alpha}^{k}\right)^{2} \tag{15}
\end{equation*}
$$

For any $T>0$, let $A_{T}$ be the superconnection on $\pi^{*} \Lambda^{*}\left(F^{*}\right)$ defined by

$$
\begin{equation*}
A_{T}=\pi^{*} \nabla^{\Lambda^{*}\left(F^{*}\right)}+c_{T}(\widehat{Y}) \tag{16}
\end{equation*}
$$

By Corollary 2.2,

$$
\begin{equation*}
\langle e(F),[M]\rangle=\left(\frac{1}{2 \pi}\right)^{2 n} \int_{\mathcal{F}} \operatorname{tr}_{s}\left[\exp \left(A_{T}^{2}\right)\right] \tag{17}
\end{equation*}
$$

We need to compute $\int_{\mathcal{F}} \operatorname{tr}_{s}\left[\exp \left(A_{T}^{2}\right)\right]$, which is independent of $T>0$. It follows from (12), (13), and (14)-(16) that

$$
\begin{align*}
A_{T}^{2} & =\left[\pi^{*} \nabla^{\Lambda^{*}\left(F^{*}\right)}, \sum_{\alpha, k} \rho_{\alpha} T^{\alpha} y_{\alpha}^{k} \widehat{d y_{\alpha}^{k}}-i_{\widehat{Y}}\right]-|Y|_{g_{T}^{F}}^{2} \\
& =\sum_{\alpha, k} T^{\alpha}\left(d \rho_{\alpha} y_{\alpha}^{k} \widehat{d y_{\alpha}^{k}}+\rho_{\alpha} d y_{\alpha}^{k} \widehat{d y_{\alpha}^{k}}\right)-\sum_{k} d y_{\alpha}^{k} \otimes i \frac{\widehat{\partial}}{\frac{\partial}{\partial y_{\alpha}^{k}}}-|Y|_{g_{T}^{F}}^{2} . \tag{18}
\end{align*}
$$

Set

$$
\begin{equation*}
\widehat{d}^{V}=\sum_{k} \widehat{d y_{\alpha}^{k}} \frac{\partial}{\partial y_{\alpha}^{k}} \tag{19}
\end{equation*}
$$

on each $\pi^{-1}\left(U_{\alpha}\right)$. Clearly, $\widehat{d}^{V}$ is well defined over $\mathcal{F}$. It follows from (15), (18), and (19) that

$$
\begin{equation*}
A_{T}^{2}=\frac{1}{2} d^{H} \widehat{d}^{V}\left(|Y|_{g_{T}^{F}}^{2}\right)+\sum_{\alpha, k} T^{\alpha} \rho_{\alpha} d y_{\alpha}^{k} \widehat{d y_{\alpha}^{k}}-\sum_{k} d y_{\alpha}^{k} \otimes i \frac{\widehat{\partial}}{\partial y_{\alpha}^{k}}-|Y|_{g_{T}^{F}}^{2} . \tag{20}
\end{equation*}
$$

Set

$$
\begin{equation*}
B_{T}^{2}=\frac{1}{2} d^{H} \widehat{d}^{V}\left(|Y|_{g_{T}^{F}}^{2}\right)-\sum_{k} d y_{\alpha}^{k} \otimes i_{\frac{\partial}{\partial y_{\alpha}^{k}}}-|Y|_{g_{T}^{F}}^{2} . \tag{21}
\end{equation*}
$$

By (20), (21), [10, Prop. 4.9], and a simple degree counting along vertical directions, one sees that

$$
\begin{equation*}
\operatorname{tr}_{s}\left[\exp \left(A_{T}^{2}\right)\right]=\operatorname{tr}_{s}\left[\exp \left(B_{T}^{2}\right)\right] \tag{22}
\end{equation*}
$$

By (19), (21), and [10, Prop. 4.9], we see that if we exchange $\widehat{d y_{\alpha}^{k}}$ and $d y_{\alpha}^{k}$, then we obtain the same supertrace of $\exp \left(B_{T}^{2}\right)$. Thus,

$$
\begin{align*}
\operatorname{tr}_{s}\left[\exp \left(B_{T}^{2}\right)\right] & =\left\{\exp \left(\frac{1}{2} d^{H} d^{V}|Y|_{g_{T}^{F}}^{2}-|Y|_{g_{T}^{F}}^{2}\right)\right\}^{(4 n)} \operatorname{tr}_{s}\left[\exp \left(-\sum_{k} \widehat{d y_{\alpha}^{k}} i_{\frac{\partial}{\partial y_{\alpha}^{k}}}\right)\right] \\
& =\left\{\exp \left(\frac{1}{2} d^{H} d^{V}|Y|_{g_{T}^{F}}^{2}-|Y|_{g_{T}^{F}}^{2}\right)\right\}^{(4 n)} . \tag{23}
\end{align*}
$$

From (17), (22), (23), and a simple transgression argument, one gets
Proposition 3.7. For any Euclidean metric $h^{F}$ on $F$, one has

$$
\langle e(F),[M]\rangle=\left(\frac{1}{2 \pi}\right)^{2 n} \int_{\mathcal{F}} \exp \left(\frac{1}{2} d^{H} d^{V}|Y|_{h^{F}}^{2}-|Y|_{h^{F}}^{2}\right) .
$$

### 3.3. Analysis outside of $B_{+}$

We continue to work with $g_{T}^{F}$.
Proposition 3.8. Every $p \in M \backslash \mathbf{B}_{+}$has an open neighborhood $V_{p} \subset M$ such that

$$
\begin{equation*}
\lim _{T \rightarrow+\infty} \int_{\mathcal{F}} f \exp \left(\frac{1}{2} d^{H} d^{V}|Y|_{g_{T}^{F}}^{2}-|Y|_{g_{T}^{F}}^{2}\right)=0 \tag{24}
\end{equation*}
$$

for any smooth function $f \in C^{\infty}(M)$ supported in $V_{p}$.
Proof. By (7) and (15), one has

$$
\frac{1}{2} d^{H} d^{V}|Y|_{g_{T}^{F}}^{2}=\sum_{\alpha, k} d \rho_{\alpha} T^{\alpha} y_{\alpha}^{k} d y_{\alpha}^{k}
$$

Thus,

$$
\begin{equation*}
\exp \left(\frac{1}{2} d^{H} d^{V}|Y|_{g_{T}^{F}}^{2}-|Y|_{g_{T}^{F}}^{2}\right)=\prod_{\alpha}\left(1+d \rho_{\alpha} T^{\alpha} \sum_{k} y_{\alpha}^{k} d y_{\alpha}^{k}\right) \exp \left(-|Y|_{g_{T}^{F}}^{2}\right) \tag{25}
\end{equation*}
$$

For simplicity, we write

$$
\begin{equation*}
h_{\alpha}=\sum_{k}\left(y_{\alpha}^{k}\right)^{2} . \tag{26}
\end{equation*}
$$

Take a $p \in M$. We assume that among $\left\{U_{\alpha}\right\}_{\alpha=1}^{N}$, there are exactly $N_{p}$ elements $U_{\alpha_{i}}$, $\alpha_{1}>\cdots>\alpha_{N_{p}}$, containing $p$. If $p$ does not lie on the boundary of any $U_{\alpha}$, then it follows from (25) and (26) that

$$
\begin{equation*}
\left\{\exp \left(\frac{1}{2} d^{H} d^{V}|Y|_{g_{T}^{F}}^{2}-|Y|_{g_{T}^{F}}^{2}\right)\right\}^{(4 n)}=\sum_{\left\{\alpha_{i_{j}}\right\}} \prod_{j=1}^{2 n}\left(d \rho_{\alpha_{i_{j}}} T^{\alpha_{i_{j}}} \frac{d h_{\alpha_{i_{j}}}}{2}\right) \exp \left(-|Y|_{g_{T}^{F}}^{2}\right) \tag{27}
\end{equation*}
$$

near $\pi^{-1}(p)$, where $\alpha_{i_{j}}$ runs through $\alpha_{i}, 1 \leq i \leq N_{p}$. Each $\alpha_{i_{j}}$ occurs at most once in a product. Moreover, by (9) and (15), one has

$$
\begin{equation*}
|Y|_{g_{T}^{F}}^{2} \geq \frac{1}{2} \rho_{\alpha_{1}}(p) T^{\alpha_{1}} \sum_{k}\left(y_{\alpha_{1}}^{k}\right)^{2} \tag{28}
\end{equation*}
$$

near $\pi^{-1}(p)$, where $\rho_{\alpha_{1}}(p)>0$. We see from (28) that there exists an open neighborhood $V_{p}$ of $p \in M$ such that

$$
\begin{equation*}
\left|\int_{\mathcal{F}} f \prod_{j=1}^{2 n}\left(d \rho_{\alpha_{i_{j}}} T^{\alpha_{i_{j}}} d h_{\alpha_{i_{j}}}\right) \exp \left(-|Y|_{g_{T}^{F}}^{2}\right)\right|=O\left(\prod_{j=1}^{2 n} T^{\alpha_{i_{j}}-\alpha_{1}}\right)=O\left(\frac{1}{T^{n(2 n-1)}}\right) \tag{29}
\end{equation*}
$$

for $T \gg 0$ for any $f \in C^{\infty}(M)$ supported in $V_{p}$. Formula (24) follows from (27) and (29).

Now let $p$ lie on the boundaries of some $U_{\alpha}$ 's. To be more precise, assume that $p$ lies on the boundaries of $\left\{U_{\beta_{i}}\right\}_{i=1}^{M_{p}}$ with $\beta_{1}>\beta_{2}>\ldots>\beta_{M_{p}}$. Then, by (25), the terms that we need to consider near $\pi^{-1}(p)$ are of the form

$$
\begin{equation*}
\prod_{h=1}^{k}\left(d \rho_{\beta_{i_{h}}} T^{\beta_{i_{h}}} d h_{\beta_{i_{h}}}\right) \prod_{j=1}^{2 n-k}\left(d \rho_{\alpha_{i_{j}}} T^{\alpha_{i_{j}}} d h_{\alpha_{i_{j}}}\right) \exp \left(-|Y|_{g_{T}^{F}}^{2}\right) \tag{30}
\end{equation*}
$$

For simplicity, we assume that $\beta_{i_{1}}>\cdots>\beta_{i_{k}}$ and $\alpha_{i_{1}}>\cdots>\alpha_{i_{2 n-k}}$. By (15), there exists a constant $c_{1}>0$ such that

$$
\begin{equation*}
|Y|_{g_{T}^{F}}^{2} \geq c_{1}\left(\sum_{j=1}^{k} \rho_{\beta_{i_{j}}} T^{\beta_{i_{j}}}+\rho_{\alpha_{1}}(p) T^{\alpha_{1}}\right) \sum_{l}\left(y_{\alpha_{1}}^{l}\right)^{2} \tag{31}
\end{equation*}
$$

near $\pi^{-1}(p)$. It follows from (31) that there exists a constant $C_{1}>0$ such that

$$
\begin{align*}
\mid \int_{\mathcal{F} / M} \prod_{h=1}^{k}\left(d \rho_{\beta_{i_{h}}} T^{\beta_{i_{h}}} d h_{\beta_{i_{h}}}\right) & \prod_{j=1}^{2 n-k}\left(d \rho_{\alpha_{i_{j}}} T^{\alpha_{i_{j}}} d h_{\alpha_{i_{j}}}\right) \exp \left(-|Y|_{g_{T}^{F}}^{2}\right) \mid \\
& \leq C_{1}\left|\frac{\prod_{h=1}^{k}\left(d \rho_{\beta_{i_{h}}} T^{\beta_{i_{h}}}\right) \prod_{j=1}^{2 n-k}\left(d \rho_{\alpha_{i_{j}}} T^{\alpha_{i_{j}}}\right)}{\left(\sum_{j=1}^{k} \rho_{\beta_{i_{j}}} T^{\beta_{i_{j}}}+\rho_{\alpha_{1}}(p) T^{\alpha_{1}}\right)^{2 n}}\right| \tag{32}
\end{align*}
$$

for $T \gg 0$ near $p \in M$, where we write $|F d \mathrm{vol}|=|F| d$ vol for a form $F d$ vol and $\mid F d$ vol $|\leq| G d$ vol $\mid$ means that $|F| \leq|G|$.

Recall that the boundaries of $U_{\beta_{i_{h}}}$ 's intersect each other transversally.
Since the function $h$ used in the definition of $\rho_{\alpha}$ in (8) is increasing, we have

$$
\begin{align*}
& \left|\frac{\prod_{h=1}^{k}\left(d \rho_{\beta_{i_{h}}} T^{\beta_{i_{h}}}\right) \prod_{j=1}^{2 n-k}\left(d \rho_{\alpha_{i_{j}}} T^{\alpha_{i_{j}}}\right)}{\left(\sum_{j=1}^{k} \rho_{\beta_{i_{j}}} T^{\beta_{i_{j}}}+\rho_{\alpha_{1}}(p) T^{\alpha_{1}}\right)^{2 n}}\right| \\
& \quad \leq\left|\prod_{h=1}^{k} \frac{d \rho_{\beta_{i_{h}}} T^{\beta_{i_{h}}}}{\rho_{\beta_{i_{h}}} T^{\beta_{i_{h}}}+\rho_{\alpha_{1}}(p) T^{\alpha_{1}}} \cdot \prod_{j=1}^{2 n-k} \frac{d \rho_{\alpha_{i_{j}}} T^{\alpha_{i_{j}}}}{\rho_{\alpha_{1}}(p) T^{\alpha_{1}}}\right|  \tag{33}\\
& \quad=\left|\left(\prod_{h=1}^{k} d \log \left(\rho_{\beta_{i_{h}}}+\rho_{\alpha_{1}}(p) T^{\alpha_{1}-\beta_{i_{h}}}\right)\right) \cdot \prod_{j=1}^{2 n-k} \frac{d \rho_{\alpha_{i_{j}}} T^{\alpha_{i_{j}}}}{\rho_{\alpha_{1}}(p) T^{\alpha_{1}}}\right|
\end{align*}
$$

It is easily seen that for $T \gg 0$ and sufficiently small $a>0$ the integration of $d \log \left(\rho_{\beta_{i_{h}}}+\rho_{\alpha_{1}}(p) T^{\alpha_{1}-\beta_{i_{h}}}\right)$ along the interval $0 \leq r_{\beta_{i_{h}}} \leq a$ gives $O(\log T)$
(respectively, $O\left(T^{-1}\right)$ ) provided that $\beta_{i_{h}}>\alpha_{1}$ (respectively, $\beta_{i_{h}}<\alpha_{1}$ ) for each $h=1, \ldots, k$. Further, if $k \leq 2 n-2$, then, in view of (29),

$$
\begin{equation*}
\prod_{j=1}^{2 n-k} \frac{T^{\alpha_{i_{j}}}}{T^{\alpha_{1}}} \leq \frac{1}{T} \quad \text { for } T \geq 1 \tag{34}
\end{equation*}
$$

One finds from (32)-(34) that if $\alpha_{1}>\beta_{i_{k}}$ or $k \leq 2 n-2$, then there exists a sufficiently small open neighborhood $V_{p}$ of $p \in M$ such that

$$
\begin{equation*}
\lim _{T \rightarrow+\infty} \int_{\mathcal{F}} f \prod_{h=1}^{k}\left(d \rho_{\beta_{i_{h}}} T^{\beta_{i_{h}}} d h_{\beta_{i_{h}}}\right) \prod_{j=1}^{2 n-k}\left(d \rho_{\alpha_{i_{j}}} T^{\alpha_{i_{j}}} d h_{\alpha_{i_{j}}}\right) \exp \left(-|Y|_{g_{T}^{F}}^{2}\right)=0 \tag{35}
\end{equation*}
$$

for any smooth function $f \in C^{\infty}(M)$ supported in $V_{p}$.
We only need to consider the case of $k=M_{p}=2 n-1$ with $\beta_{2 n-1}>\alpha_{1}$ in (30) and the case of $M_{p}=2 n$. In the former case, one still obtains (35) by (33) if $\alpha_{i_{1}}<\alpha_{1}$, and it only remains to consider the term

$$
\begin{align*}
& \prod_{h=1}^{2 n-1}\left(d \rho_{\beta_{h}} T^{\beta_{h}} d h_{\beta_{h}}\right)\left(d \rho_{\alpha_{1}} T^{\alpha_{1}} d h_{\alpha_{1}}\right) \exp \left(-|Y|_{g_{T}^{F}}^{2}\right) \\
& \quad=d^{V}\left(\prod_{h=1}^{2 n-1}\left(d \rho_{\beta_{h}} T^{\beta_{h}} d h_{\beta_{h}}\right) \rho_{\alpha_{1}}^{-1} d \rho_{\alpha_{1}} \exp \left(-|Y|_{g_{T}^{F}}^{2}\right)\right)  \tag{36}\\
& \quad-\prod_{h=1}^{2 n-1}\left(d \rho_{\beta_{h}} T^{\beta_{h}} d h_{\beta_{h}}\right) \rho_{\alpha_{1}}^{-1} d \rho_{\alpha_{1}}\left(\sum_{i=2}^{N_{p}} \rho_{\alpha_{i}} T^{\alpha_{i}} d h_{\alpha_{i}}\right) \exp \left(-|Y|_{g_{T}^{F}}^{2}\right),
\end{align*}
$$

from which we still obtain (35) via (33) using the Stokes formula.
For the case of $M_{p}=2 n$, one has $\rho_{\alpha_{j}}=1, j=1, \ldots, N_{p}$, near $p \in M$ by Remark 3.1. Thus, one may assume that $k=2 n$. Since $\beta_{2 n}(p)<\alpha_{1}(p)$, we see that (35) follows from (33).

Thus, (35), which implies (24), holds for $p \notin \mathbf{B}_{+}$.

### 3.4. Proof of Proposition 3.2

Now assume that $p \in \mathbf{B}_{+}$. Recall that $p$ is a point of intersection of $\left\{\partial U_{\beta_{i}}\right\}_{i=1}^{2 n}$, $\beta_{1}>\ldots>\beta_{2 n}$, and lies in the open coordinate charts $\left\{U_{\alpha_{j}}\right\}_{j=1}^{N_{p}}, \alpha_{1}>\cdots>\alpha_{N_{p}}$. Moreover, $\beta_{2 n}>\alpha_{1}$. For brevity, set

$$
\begin{equation*}
|Y|_{T}^{2}=\sum_{i=1}^{2 n} \rho_{\beta_{i}} T^{\beta_{i}} h_{\beta_{i}}+T^{\alpha_{1}} h_{\alpha_{1}} \tag{37}
\end{equation*}
$$

on $\pi^{-1}\left(W_{p}\right)$.
Let $\phi, 0 \leq \phi \leq 1$, be a smooth function on $M$ such that $\operatorname{Supp}(\phi) \subset V_{p}$, where $V_{p} \subset W_{p}$ is a sufficiently small open neighborhood of $p$, and $\phi=1$ near
$p \in M$. We only need to prove that there exists a limit

$$
\lim _{T \rightarrow+\infty} \int_{\mathcal{F}} \phi \prod_{j=1}^{2 n}\left(d \rho_{\beta_{j}} T^{\beta_{j}} d h_{\beta_{j}}\right) \exp \left(-|Y|_{T}^{2}\right)
$$

For any $T>0$, set

$$
\begin{equation*}
\gamma_{T}=\frac{\partial}{\partial T} \int_{\mathcal{F}} \phi \prod_{j=1}^{2 n}\left(d \rho_{\beta_{j}} T^{\beta_{j}} d h_{\beta_{j}}\right) \exp \left(-|Y|_{T}^{2}\right) \tag{38}
\end{equation*}
$$

Lemma 3.9. For $T \gg 0$, one has

$$
\begin{equation*}
\gamma_{T}=O\left(\frac{(\log T)^{2 n-1}}{T^{2}}\right) \tag{39}
\end{equation*}
$$

Proof. Set $i_{Y}=\sum_{k} y_{\alpha}^{k} i_{\frac{\partial}{\partial y_{\alpha}^{k}}}$ on each $\pi^{-1}\left(U_{\alpha}\right)$; this is independent of $\alpha$. Then

$$
\begin{equation*}
\left(d^{H}-i_{Y}\right)^{2}=0 . \tag{40}
\end{equation*}
$$

One can verify that

$$
\begin{equation*}
\left(d^{H}-i_{Y}\right)\left(\rho_{\alpha} d h_{\alpha}\right)=d \rho_{\alpha} d h_{\alpha}-2 \rho_{\alpha} h_{\alpha} \tag{41}
\end{equation*}
$$

for $1 \leq \alpha \leq N$. By (37), (38), (40), (41), and the Stokes formula,

$$
\begin{align*}
& \gamma_{T}= 2^{2 n} \frac{\partial}{\partial T} \int_{\mathcal{F}} \phi \prod_{i=1}^{2 n}\left(e^{\left(d^{H}-i_{Y}\right)\left(\rho_{\beta_{i}} T^{\beta_{i}} d h_{\beta_{i}} / 2\right)}\right) e^{\left(d^{H}-i_{Y}\right) T^{\alpha_{1}} d h_{\alpha_{1}} / 2} \\
&= \frac{2^{2 n-1}}{T} \int_{\mathcal{F}} \phi\left(d^{H}-i_{Y}\right)\left(\left(\sum_{i=1}^{2 n} \beta_{i} \rho_{\beta_{i}} T^{\beta_{i}} d h_{\beta_{i}}+\alpha_{1} T^{\alpha_{1}} d h_{\alpha_{1}}\right)\right. \\
&\left.\cdot\left(\prod_{i=1}^{2 n} e^{\left(d^{H}-i_{Y}\right)\left(\rho_{\beta_{i}} T^{\beta_{i}} d h_{\beta_{i}} / 2\right)}\right) e^{\left(d^{H}-i_{Y}\right) T^{\alpha_{1}} d h_{\alpha_{1}} / 2}\right) \\
&=-\frac{1}{T} \int_{\mathcal{F}} d \phi\left(\sum_{i=1}^{2 n} \beta_{i} \rho_{\beta_{i}} T^{\beta_{i}} d h_{\beta_{i}} \prod_{j \neq i}\left(d \rho_{\beta_{j}} T^{\beta_{j}} d h_{\beta_{j}}\right)\right) \exp \left(-|Y|_{T}^{2}\right)  \tag{42}\\
& \quad-\frac{\alpha_{1}}{T} \int_{\mathcal{F}} d \phi T^{\alpha_{1}} d h_{\alpha_{1}}\left(\sum_{i=1}^{2 n} \prod_{j \neq i}\left(d \rho_{\beta_{j}} T^{\beta_{j}} d h_{\beta_{j}}\right)\right) \exp \left(-|Y|_{T}^{2}\right) \\
&= \frac{1}{T} \int_{\mathcal{F}} d \phi T^{\alpha_{1}} d h_{\alpha_{1}}\left(\sum_{i=1}^{2 n} \beta_{i} \prod_{j \neq i}\left(d \rho_{\beta_{j}} T^{\beta_{j}} d h_{\beta_{j}}\right)\right) \exp \left(-|Y|_{T}^{2}\right) \\
& \quad-\frac{\alpha_{1}}{T} \int_{\mathcal{F}} d \phi T^{\alpha_{1}} d h_{\alpha_{1}}\left(\sum_{i=1}^{2 n} \prod_{j \neq i}\left(d \rho_{\beta_{j}} T^{\beta_{j}} d h_{\beta_{j}}\right)\right) \exp \left(-|Y|_{T}^{2}\right),
\end{align*}
$$

where the last equality follows from a vertical transgression argument (cf. (36)).
For any $q \in \operatorname{Supp}(d \phi)$, either one of the numbers $\rho_{\beta_{i}}(q)$ is positive, or one of the $\rho_{\beta_{i}}$ 's vanishes near $q$. In the former case, since $\beta_{i}>\alpha_{1}$, we can proceed as
in (33) and prove that there exists a small open neighborhood $V_{q} \subset W_{p}$ of $q$ such that

$$
\begin{align*}
& \int_{\mathcal{F}} f d \phi T^{\alpha_{1}} d h_{\alpha_{1}}\left(\sum_{i=1}^{2 n} \beta_{i} \prod_{j \neq i}\left(d \rho_{\beta_{j}} T^{\beta_{j}} d h_{\beta_{j}}\right)\right) \exp \left(-|Y|_{T}^{2}\right)  \tag{43}\\
& -\alpha_{1} \int_{\mathcal{F}} f d \phi T^{\alpha_{1}} d h_{\alpha_{1}}\left(\sum_{i=1}^{2 n} \prod_{j \neq i}\left(d \rho_{\beta_{j}} T^{\beta_{j}} d h_{\beta_{j}}\right)\right) \exp \left(-|Y|_{T}^{2}\right)=O\left(\frac{(\log T)^{2 n-1}}{T}\right)
\end{align*}
$$

for $T \gg 1$ for any $f \in C^{\infty}(M)$ with $\operatorname{Supp}(f) \subset V_{q}$. In the later case, by an easy vertical transgression argument,

$$
\begin{align*}
& \int_{\mathcal{F} / M} d \phi T^{\alpha_{1}} d h_{\alpha_{1}}\left(\sum_{i=1}^{2 n} \beta_{i} \prod_{j \neq i}\left(d \rho_{\beta_{j}} T^{\beta_{j}} d h_{\beta_{j}}\right)\right) \exp \left(-|Y|_{T}^{2}\right)  \tag{44}\\
&-\alpha_{1} \int_{\mathcal{F} / M} d \phi T^{\alpha_{1}} d h_{\alpha_{1}}\left(\sum_{i=1}^{2 n} \prod_{j \neq i}\left(d \rho_{\beta_{j}} T^{\beta_{j}} d h_{\beta_{j}}\right)\right) \exp \left(-|Y|_{T}^{2}\right)=0
\end{align*}
$$

near $q \in M$. Now one obtains (39) by combining (42)-(44) with a simple partition of unity argument.

From (38) and Lemma 3.9, one sees that there exists a limit

$$
\begin{aligned}
& \lim _{T \rightarrow+\infty} \int_{\mathcal{F}} \phi \prod_{j=1}^{2 n}\left(d \rho_{\beta_{j}} T^{\beta_{j}} d h_{\beta_{j}}\right) \exp \left(-|Y|_{T}^{2}\right) \\
&=\int_{\mathcal{F}} \phi \prod_{j=1}^{2 n}\left(d \rho_{\beta_{j}} d h_{\beta_{j}}\right) \exp \left(-|Y|_{T=1}^{2}\right)+\lim _{T \rightarrow+\infty} \int_{1}^{T} \gamma_{t} d t
\end{aligned}
$$

which completes the proof of Proposition 3.2.

### 3.5. Proof of Theorem 3.4

First, we still assume that $p \in \mathbf{B}_{+}$. Recall that $\rho_{\alpha_{j}}=1,1 \leq j \leq N_{p}$, on $W_{p}$ by Remark 3.1. Further,

$$
\begin{equation*}
|Y|_{g_{T}^{F}}^{2}=\sum_{i=1}^{2 n} \rho_{\beta_{i}} T^{\beta_{i}} h_{\beta_{i}}+\sum_{i=1}^{N_{p}} T^{\alpha_{i}} h_{\alpha_{i}} \tag{45}
\end{equation*}
$$

near $\pi^{-1}(p)$. We only need to consider the term

$$
\lim _{T \rightarrow+\infty} \int_{\mathcal{F}} \phi \prod_{i=1}^{2 n}\left(d \rho_{\beta_{i}} T^{\beta_{i}} d h_{\beta_{i}}\right) \exp \left(-|Y|_{g_{T}^{F}}^{2}\right)
$$

which is examined in the following lemma.

Lemma 3.10. The following identity holds:

$$
\begin{align*}
\lim _{T \rightarrow+\infty} \int_{\mathcal{F}} & \phi \prod_{i=1}^{2 n}\left(d \rho_{\beta_{i}} T^{\beta_{i}} d h_{\beta_{i}}\right) \exp \left(-|Y|_{g_{T}^{F}}^{2}\right) \\
& =\lim _{T \rightarrow+\infty} \int_{\mathcal{F}} \phi \prod_{i=1}^{2 n}\left(d \rho_{\beta_{i}} T^{\beta_{i}} d h_{\beta_{i}}\right) \exp \left(-|Y|_{T}^{2}\right) \tag{46}
\end{align*}
$$

Proof. It follows from (37), (40), (41), and (45) that

$$
\begin{gather*}
\frac{\phi}{2^{2 n}} \prod_{i=1}^{2 n}\left(d \rho_{\beta_{i}} T^{\beta_{i}} d h_{\beta_{i}}\right) \exp \left(-|Y|_{g_{T}^{F}}^{2}\right)-\frac{\phi}{2^{2 n}} \prod_{i=1}^{2 n}\left(d \rho_{\beta_{i}} T^{\beta_{i}} d h_{\beta_{i}}\right) \exp \left(-|Y|_{T}^{2}\right) \\
=\left\{\phi \prod_{i=1}^{2 n}\left(e^{\left(d^{H}-i_{Y}\right)\left(\rho_{\beta_{i}} T^{\beta_{i}} d h_{\beta_{i}} / 2\right)}\right) e^{-T^{\alpha_{1}} h_{\alpha_{1}}}\left(e^{\left(d^{H}-i_{Y}\right) \sum_{i=2}^{N_{p}} T^{\alpha_{i}} d h_{\alpha_{i}} / 2}-1\right)\right\}^{(4 n)} \\
=\left\{( d ^ { H } - i _ { Y } ) \left(\phi \prod_{i=1}^{2 n}\left(e^{\left(d^{H}-i_{Y}\right)\left(\rho_{\beta_{i}} T^{\beta_{i}} d h_{\beta_{i}} / 2\right)}\right) e^{-T^{\alpha_{1}} h_{\alpha_{1}}}\left(\sum_{i=2}^{N_{p}} T^{\alpha_{i}} d h_{\alpha_{i}} / 2\right)\right.\right. \\
\left.\cdot \frac{e^{-\sum_{i=2}^{N_{p}} T^{\alpha_{i}} h_{\alpha_{i}}}-1}{-\sum_{i=2}^{N_{p}} T^{\alpha_{i}} h_{\alpha_{i}}}\right)  \tag{47}\\
-d \phi \prod_{i=1}^{2 n}\left(e^{\left(d^{H}-i_{Y}\right)\left(\rho_{\beta_{i}} T^{\beta_{i}} d h_{\beta_{i}} / 2\right)}\right) e^{-T^{\alpha_{1}} h_{\alpha_{1}}}\left(\sum_{i=2}^{N_{p}} T^{\alpha_{i}} d h_{\alpha_{i}} / 2\right) \\
\left.\cdot \frac{e^{-\sum_{i=2}^{N_{p}} T^{\alpha_{i}} h_{\alpha_{i}}}-1}{-\sum_{i=2}^{N_{p}} T^{\alpha_{i}} h_{\alpha_{i}}}\right\}^{(4 n)}
\end{gather*}
$$

From (37), (41), (47), and the Stokes formula, one obtains

$$
\begin{gather*}
\int_{\mathcal{F}} \phi \prod_{h=1}^{2 n}\left(d \rho_{\beta_{h}} T^{\beta_{h}} d h_{\beta_{h}}\right) \exp \left(-|Y|_{g_{T}^{F}}^{2}\right)-\int_{\mathcal{F}} \phi \prod_{h=1}^{2 n}\left(d \rho_{\beta_{h}} T^{\beta_{h}} d h_{\beta_{h}}\right) \exp \left(-|Y|_{T}^{2}\right) \\
=\int_{\mathcal{F}} d \phi\left(\sum_{h=1}^{2 n} \prod_{i \neq h}\left(d \rho_{\beta_{i}} T^{\beta_{i}} d h_{\beta_{i}}\right)\right)\left(\sum_{i=2}^{N_{p}} T^{\alpha_{i}} d h_{\alpha_{i}}\right)  \tag{48}\\
\cdot \frac{e^{-\sum_{i=2}^{N_{p}} T^{\alpha_{i}} h_{\alpha_{i}}}-1}{\sum_{i=2}^{N_{p}} T^{\alpha_{i}} h_{\alpha_{i}}} \exp \left(-|Y|_{T}^{2}\right)
\end{gather*}
$$

Recall that $\alpha_{1}>\alpha_{i}$ for $i \geq 2$. Further,

$$
\begin{equation*}
0 \leq \frac{1-e^{-t}}{t} \leq 1 \tag{49}
\end{equation*}
$$

for any $t \geq 0$. We derive (46) from (37), (48), and (49), arguing as in (33).
Returning to the proof of Theorem 3.4, we take a finite selection of $V_{p}$ 's in Propositions 3.2 and 3.8 so that they form an open cover of $M$. Let $\left\{f_{V_{p}}\right\}$ be a partition
of unity subordinate to this cover. We assume that each $p \in \mathbf{B}_{+}$is covered by only one $V_{p}$ on which $f_{V_{p}}=1$ near $p$. Since $\mathbf{B}_{+}$consists of finitely many points, the existence of such a cover and partition of unity is clear.

Now Theorem 3.4 follows from Propositions 3.2, 3.7, and 3.8 and Lemma 3.10:

$$
\langle e(F),[M]\rangle=\left(\frac{1}{2 \pi}\right)^{2 n} \sum_{T \rightarrow+\infty} \lim _{\mathcal{F}} f_{V_{p}} \exp \left(\frac{1}{2} d^{H} d^{V}|Y|_{g_{T}^{F}}^{2}-|Y|_{g_{T}^{F}}^{2}\right)=\sum_{q \in \mathbf{B}_{+}} \nu_{q} .
$$

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[^0]:    ${ }^{1}$ The ordering of $U_{\alpha}$ 's, while arbitrary, will play an important role in what follows.

[^1]:    ${ }^{2}$ It is easily seen that there always exists a transversal open cover.

